

Erratum: Spin of Dirac's relativistic membrane
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(i) Equations (3.6)–(3.8) should be replaced by

$$\begin{aligned}
 \mathbf{I}T^\dagger T = & (p^2 + \mu^2 + q^2 + k^2 f^2 - \hbar \nabla \cdot \mathbf{q}) \mathbf{I} + [\hbar \nabla \cdot \mathbf{f} - 2(\mathbf{q} \cdot \mathbf{f})] \kappa \beta - i[\hbar \boldsymbol{\alpha} \cdot \nabla \mu - 2\mu(\boldsymbol{\alpha} \cdot \mathbf{q})] \beta \\
 & + i[\{q_x, p_y\} - \{p_x, q_y\} - \kappa(\{f_x, p_y\} - \{p_x, f_y\})] \beta \boldsymbol{\alpha}_x \boldsymbol{\alpha}_y \\
 & + i[\{q_x, p_z\} - \{p_x, q_z\} - \kappa(\{f_x, p_z\} - \{p_x, f_z\})] \beta \boldsymbol{\alpha}_x \boldsymbol{\alpha}_z \\
 & + i[\{q_y, p_z\} - \{p_y, q_z\} - \kappa(\{f_y, p_z\} - \{p_y, f_z\})] \beta \boldsymbol{\alpha}_y \boldsymbol{\alpha}_z,
 \end{aligned} \tag{3.6}$$

since $\mathbf{p} = -i\hbar \nabla$ and where κ is the eigenvalue of K and also assuming that K is a linear Hermitian operator. The linearization we look for requires $T^2 = p^2 + \mu^2$, and thus we conclude from the above matrix equation that we must impose $q_x = f_x$, $q_y = f_y$, $q_z = f_z$, i.e., $\mathbf{q} = \mathbf{f}$ and also $\kappa = \pm 1$. Otherwise, in order that the third term of Eq. (3.6) vanishes we must have $q_x = \hbar x/r^2$, $q_y = \hbar y/r^2$ and $q_z = \hbar z/r^2$, where $r^2 = x^2 + y^2 + z^2$. Hence, owing to the above results the linearized form of the Hamiltonian operator is

$$H = \boldsymbol{\alpha} \cdot \mathbf{P} + i(\boldsymbol{\alpha} \cdot \mathbf{q}) \beta K + \beta \mu + \mathbf{I}V, \tag{3.7}$$

according to Eq. (3.4), where $\mathbf{P} = \mathbf{p} - i\mathbf{q}$, $p_j = -i\hbar \partial/\partial x^j$, $q^j = \hbar x^j/r^2$, and $|\mathbf{q}| = \hbar/r$.

(ii) The commutation rules (5.1) should be

$$[x^j, P_k] = i\hbar \delta_k^j, \quad [P_j, P_k] = 0, \quad [\Omega_j, P_k] = i\hbar \varepsilon_{jkl} P_l, \quad [\Omega_j, q_k] = i\hbar \varepsilon_{jkl} q_l, \tag{5.1}$$

where Ω_z does not commute with f_x and f_y . Since in this case $\Omega_x = \Omega_y = 0$ and assuming that Ω_z commutes with the operator K , then with the Hamiltonian (3.7) we obtain for the Ω_z component the result given in Eq. (5.2).

(iii) Below Eq. (5.11) we have $\mathbf{P}_r = 1/r(\mathbf{r} \cdot \mathbf{P})$ and the matrices $\boldsymbol{\sigma}$ that appear in Eqs. (5.11), (5.13)–(5.15) are indeed the 2×2 Pauli spin matrices.

The above corrections do not alter the results derived in this paper.